

# A Tanaka formula for the derivative of self-intersection local time for fBm

Paul Jung (University of Alabama Birmingham)  
Greg Markowsky (Monash University, Melbourne)

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# Self-intersection local time

Brownian **local time** is formally defined by

$$L(t, x) := \int_0^t \delta(x - B_s) ds.$$

Two related processes are the **Intersection Local Time**:

$$\int_0^t \int_0^t \delta(B_s - \tilde{B}_r) dr ds$$

and **Self-intersection Local Time (SLT)**:

$$\int_0^t \int_0^s \delta(B_s - B_r) dr ds$$

SLT is used in physics to study polymers, the polaron, and QFT [see X. Chen (2008, 2010)].

# The Derivative of Self-intersection Local Time (DSLTL)

Our interested is a formal **Derivative of SLT**:

$$\alpha'(t) := -\frac{1}{2} \int_0^t \int_0^s \delta'(B_s - B_r) dr ds$$

- Introduced by Rogers/Walsh (90, 91a, 91b) for studying stochastic area integrals w.r.t. local time.
- Rosen (2005) and Markowsky (2008) showed a Tanaka formula:

$$\alpha'(t) = \int_0^t L_s^{B_s} dB_s - \frac{1}{2} \int_0^t \text{sgn}(B_t - B_r) dr.$$

- Y. Hu and D. Nualart (2009,2010) used  $\alpha'(t)$  to show CLTs for the  $L^{2,3}$  moduli of continuity for Brownian local time.
- $\alpha'(t)$  has finite nonzero 4/3-variation [Rogers/Walsh (91b) and Hu/Nualart/Song (2012)].

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# The DSLT of fBm

Recall that fBm with  $H \in (0, 1)$  is the centered Gaussian process with covariance

$$\frac{1}{2}(s^{2H} + t^{2H} - |t - s|^{2H}).$$

For  $H < 2/3$ , Yan, Yang, Lu (2008) introduce a version of  $\alpha'(t)$  connected with stochastic area integrals w.r.t. local times. We modify their definition to

$$\alpha'(t) := -H \int_0^t \int_0^s \delta'(B_s^H - B_r^H)(s - r)^{2H-1} dr ds.$$

Existence of  $\alpha'(t)$

For  $H < 2/3$ ,  $\alpha'(t)$  exists as a limit in  $L^2(\Omega)$  of

$$\alpha'_\varepsilon(t) := -H \int_0^t \int_0^s f'_\varepsilon(B_s^H - B_r^H)(s - r)^{2H-1} dr ds.$$

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## Theorem

For  $H < 2/3$ ,  $\alpha'(t)$  exists in  $L^2(\Omega)$ . Its Wiener chaos is

$$\alpha'(t) = \sum_{m=1}^{\infty} I_{2m-1}(g(2m-1, t))$$

where

$$\begin{aligned} & g(2m-1, t; v_1, \dots, v_{2m-1}) \\ = & \frac{(-1)^m}{(m-1)!2^{m-1}\sqrt{2\pi}} \int_0^t \int_0^s \frac{\prod_{j=1}^{2m-1} M_H 1_{[r,s]}(v_j) dr ds}{(s-r)^{H(2m-1)+1}} \end{aligned}$$

and

$$\langle M_H 1_{[0,s]}, M_H 1_{[0,t]} \rangle_{L^2(\mathbb{R})} = \frac{1}{2}(t^{2H} + s^{2H} - |t-s|^{2H}).$$



Apply Stroock's formula to  $\alpha'_\epsilon(t)$ :

$$\frac{1}{n!} \int_0^t \int_0^s (s-r)^{2H-1} \mathbf{E}[D^n f'_\epsilon(B_s^H - B_r^H)] dr ds$$

and then use the following:

## Lemma (Nualart and Vives (92))

Let  $F_\epsilon$  be a family of  $L^2(\Omega)$  random variables with chaos expansions  $F_\epsilon = \sum_{n=0}^{\infty} I_n(f_n^\epsilon)$ . If for all  $n$ ,  $f_n^\epsilon$  converges in  $\mathcal{H}^{\otimes n}$  to  $f_n$ , and if

$$\sum_{n=0}^{\infty} \sup_{\epsilon} \mathbf{E}[I_n(f_n^\epsilon)^2] = \sum_{n=0}^{\infty} \sup_{\epsilon} \{n! \|f_n^\epsilon\|_{\mathcal{H}^{\otimes n}}^2\} < \infty,$$

then  $F_\epsilon$  converges in  $L^2(\Omega)$  to  $F = \sum_{n=0}^{\infty} I_n(f_n)$ .

# A Tanaka formula for the DSLT of fBm

## Theorem

For  $H < 2/3$ , the following holds in  $L^2(\Omega)$  for all  $t$ :

$$H \alpha'(t) = \int_0^t L_s^{B_s^H} dB_s^H - \frac{1}{2} \int_0^t \text{sgn}(B_t^H - B_r^H) dr .$$

Formally apply a fractional Itô formula [e.g., Bender (03)] to  $B_s^H - B_r^H$  with  $s$  going from  $r$  to  $t$ :

$$\begin{aligned} & 1_{[(0,\infty)}(B_t^H - B_r^H) - 1_{[(0,\infty)}(0) \\ &= \int_r^t \delta(B_s^H - B_r^H) dB_s^H + H \int_r^t \delta'(B_s^H - B_r^H)(s - r)^{2H-1} ds \end{aligned}$$

Then integrate  $dr$  from 0 to  $t$  and apply Fubini:

$$\begin{aligned} & \int_0^t \frac{1}{2} \text{sgn}(B_t^H - B_r^H) dr \\ &= \int_0^t L_s^{B_s^H} dB_s^H + H \int_0^t \int_0^s \delta'(B_s^H - B_r^H)(s - r)^{2H-1} dr ds . \end{aligned}$$

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More rigorously, the first term on the LHS is the limit of

$$\int_0^t \int_r^t f_\varepsilon(B_s^H - B_r^H) dB_s^H dr$$

To justify Fubini, use the chaos expansion:

$$f_\varepsilon(B_s^H - B_r^H) = \sum_{n \geq 0} I_n \left( \frac{1}{n!} E \left[ \left( \frac{d^n}{dx^n} f_\varepsilon \right) (B_s^H - B_r^H) \right] (M_H 1_{[r,s]})^{\otimes n} \right).$$

We refine this to a **Hermite chaos expansion** using (tensor products of) Hermite functions as a basis of  $\mathbf{H}^{\otimes n}$ :

$$c_\beta(s, r) = \frac{1}{n!} E \left[ \left( \frac{d^n}{dx^n} f_\varepsilon \right) (B_s^H - B_r^H - y) \right] \langle \xi^{\odot \beta}, (M_H 1_{[r,s]})^{\otimes n} \rangle_{\mathbf{H}^{\otimes n}}.$$

If  $\sum \beta! c_\beta < \infty$  then we are in  $L^2(\Omega)$ . For more stringent summability, we get **Hida test functions** (and thus distributions).

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# A Fubini theorem

An extension of Fubini theorems in Cheridito/Nualart (2005) and Mishura (2008) to integrals of Hida distributions.

## Lemma (Fubini-Tonelli theorem)

Let

$$F_{s,r} = \sum_{\beta \in \Lambda} c_{\beta}(s,r) \mathcal{H}_{\beta}$$

be an  $(\mathcal{S})^*$ -valued process indexed by  $(s,r) \in \mathbb{R} \times [0, t]$ . If, for each  $(\beta, k)$  pair,  $c_{\beta}(s,r) M_{H\xi_k}(s)$  is bounded above or below by an  $L^1([r, t] \times [0, t])$  function, then

$$\int_0^t \int_r^t F_{s,r}(\omega) dB_s^H dr = \int_0^t \left( \int_0^s F_{s,r}(\omega) dr \right) dB_s^H. \quad (1)$$

The equality in (1) is in the sense that if one side is in  $(\mathcal{S})^*$ , then the other is as well, and they are equal.



# Back to existence of $\alpha'(t)$

Set

$$\mathcal{D}_t := \{0 \leq r \leq s \leq t\}$$

$$\lambda := \text{Var}(B_s^H - B_r^H)$$

$$\rho := \text{Var}(B_{s'}^H - B_{r'}^H)$$

$$\mu := \text{Cov}(B_s^H - B_r^H, B_{s'}^H - B_{r'}^H)$$

Existence in  $L^2(\Omega)$  is implied by

$$\int_{\mathcal{D}_t^2} \frac{\mu(s-r)^{2H-1}(s'-r')^{2H-1}}{(\lambda\rho - \mu^2)^{3/2}} dr dr' ds ds' < \infty.$$

Comes from looking at  $\mathbf{E}[(\alpha'_\epsilon(t))^2]$ .

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## Second moment bounds from Hu (2001)

$$\lambda = \text{Var}(B_s^H - B_r^H), \quad \rho = \text{Var}(B_{s'}^H - B_{r'}^H), \quad \mu = \text{Cov}(B_s^H - B_r^H, B_{s'}^H - B_{r'}^H)$$

Lemma 3.1 of Hu (2001) asserts

(i) For  $r < r' < s < s'$  and  $a = r' - r$ ,  $b = s - r'$ ,  $c = s' - s$ ,

$$\lambda\rho - \mu^2 \geq C \left( (a+b)^{2H} c^{2H} + a^{2H} (b+c)^{2H} \right).$$

(ii) For  $r < r' < s' < s$  and  $a = r' - r$ ,  $b = s' - r'$ ,  $c = s - s'$ ,

$$\lambda\rho - \mu^2 \geq C b^{2H} (a+b+c)^{2H}.$$

(iii) For  $r < s < r' < s'$  and  $a = s - r$ ,  $b = r' - s$ ,  $c = s' - r'$ ,

$$\lambda\rho - \mu^2 \geq C (a^{2H} c^{2H}).$$

## A modification of (ii)

In every instance we have seen, the following suffices:

(ii') For  $r < r' < s' < s$  and  $a = r' - r$ ,  $b = s' - r'$ ,  $c = s - s'$ ,

$$\lambda\rho - \mu^2 \geq Cb^{2H}(a^{2H} + c^{2H}).$$

## Conjecture

- The critical parameter is  $H_c = 2/3$ . At  $H_c$ ,  $\frac{1}{\log(1/\varepsilon)^\gamma} \alpha'_\varepsilon(t)$  converges in distribution to a normal law for some  $\gamma > 0$ .
- For  $H > H_c$ ,  $\varepsilon^{-\gamma(H)} \alpha'_\varepsilon(t)$  converges in distribution to a normal law for some function  $\gamma(H) > 0$  which goes to 0 at  $H = 2/3$ .

- We have shown related results for

$$\tilde{\alpha}'(t) := -H \int_0^t \int_0^s \delta'(B_s^H - B_r^H) dr ds.$$

- This would mirror the behavior of SLT shown in [Y. Hu and D. Nualart \(2005\)](#). To our knowledge, no such CLT has been proved even for SLT in two dimensions at  $H_c = 3/4$ .
- This would validate Rosen's statement that "The (DSL)T of Brownian motion) in  $\mathbb{R}^1$ , in a certain sense, is even more singular than self-intersection local time in  $\mathbb{R}^2$ ."

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# A generalization to $\alpha'_t(y)$

The above generalizes to

$$\alpha'_t(y) := -H \int_0^t \int_0^s \delta'(B_s^H - B_r^H - y)(s - r)^{2H-1} dr ds.$$

## Theorem

For  $H < 2/3$ , the following holds in  $L^2(\Omega)$  for all  $y$  and  $t$ :

$$H \alpha'_t(y) + \frac{1}{2} \operatorname{sgn}(y)t = \int_0^t L_s^{B_s^H - y} dB_s^H - \frac{1}{2} \int_0^t \operatorname{sgn}(B_t^H - B_r^H - y) dr .$$



Thanks for your attention!